Photoinduced structural dynamics

investigated by pump-probe diffraction

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Crystallography and large-scale facilities 2024

Photoinduced structural dynamics

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I. Scientific context

- II. Pump-probe diffraction
 - Pump-probe scheme
 - Time resolution, synchronization
 - Short X-ray pulse sources
 - Specific geometrical constraints

III. Tracking structural dynamics with pump-probe diffraction

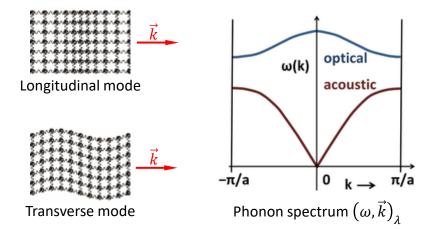
- Structural information brought by diffraction
- Coherent phonons
- Photoinduced phase transitions
- Photoinduced strain and temperature elevation

Photoinduced structural dynamics investigated by pump-probe diffraction Part I. Scientific context

1- Crystals at thermodynamic equilibrium

Atomic displacements : sum of normal modes

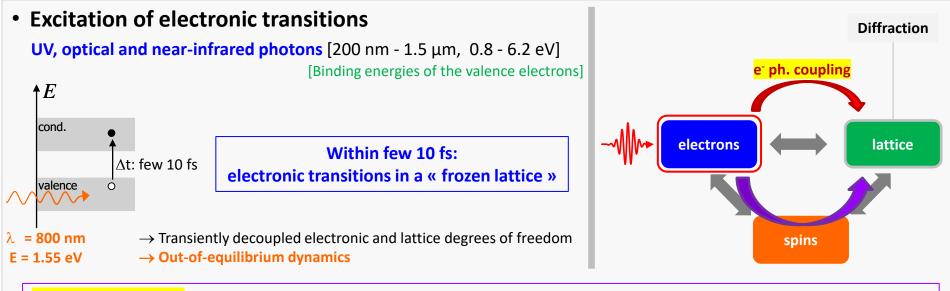
$$\vec{u_n}(\vec{r},t) = \sum_{\lambda, \|\vec{k}\|} u_n(\lambda, \vec{k}) \vec{e}_{\lambda, \vec{k}} e^{i[\omega(\lambda, \vec{k})t - \vec{k}.\vec{r}]}$$



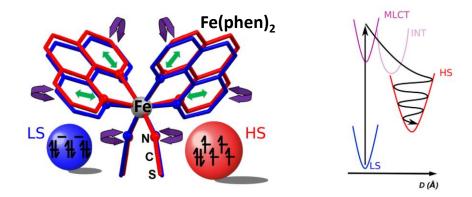
→ Experiments in the <u>frequency domain</u>: inelastic neutron scattering, Raman scattering...

Photoinduced structural dynamics investigated by pump-probe diffraction Part I. Scientific context

2- Photoinduced, out-of-equilibrium structural dynamics



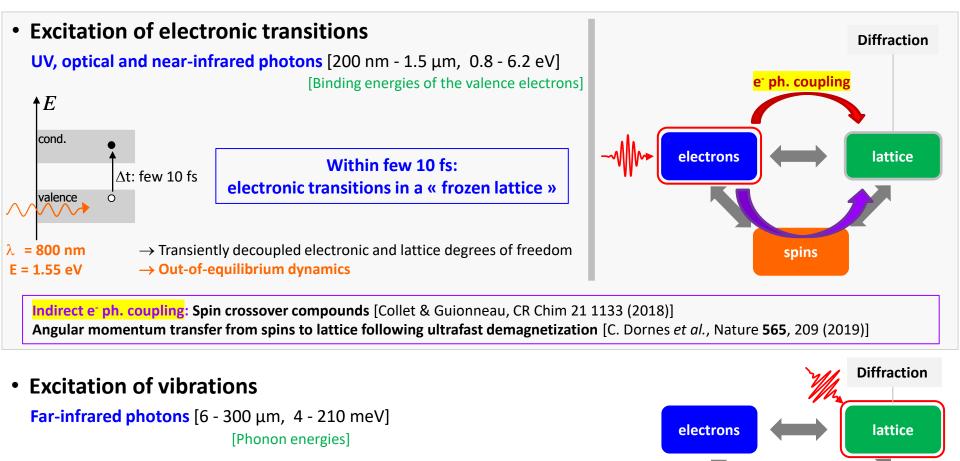
Indirect e ph. coupling: Spin crossover compounds [Collet & Guionneau, CR Chim 21 1133 (2018)] Angular momentum transfer from spins to lattice following ultrafast demagnetization [C. Dornes et al., Nature 565, 209 (2019)]



→ Experiments in the time domain: time-resolved pump-probe diffraction

Photoinduced structural dynamics investigated by pump-probe diffraction Part I. Scientific context

2- Photoinduced, out-of-equilibrium structural dynamics



 \rightarrow Direct excitation of a phonon mode

 \rightarrow Indirect excitation of phonon modes via anharmonic couplings ("nonlinear phononics")

Mankowsky et al., Rep. Prog. Phys. 79, 064503 (2016)

→ Experiments in the <u>time domain</u>: time-resolved pump-probe diffraction

spins

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Scientific context

II. Pump-probe diffraction

- Pump-probe scheme
- Time resolution, synchronization
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- Specific geometrical constraints

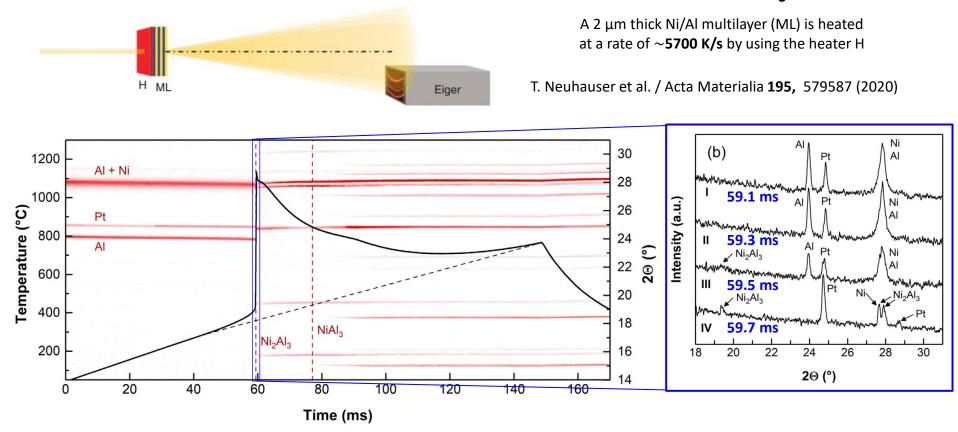
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0- Real time experiments

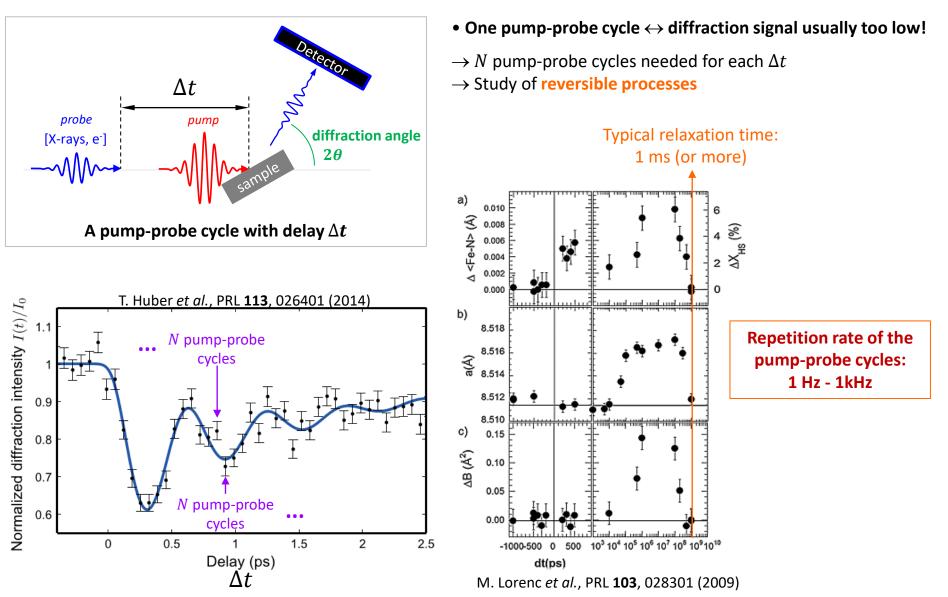
• Example: study of the exothermic solid-state transformation $[Ni+A] \rightarrow NiAl_3]$



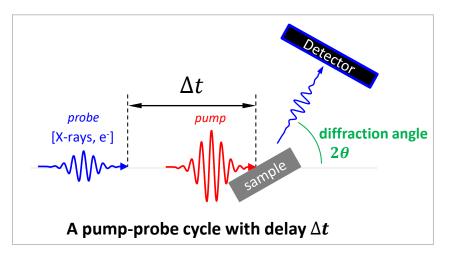
Several factors limit the time resolution in « real time » experiments:

- Limited detector frame rate
- Limited X-ray flux
- > Need to scan sample angles to retrieve the relevant information

1- Pump-probe scheme

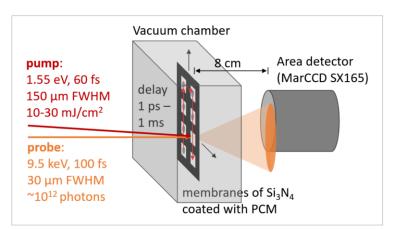


1- Pump-probe scheme

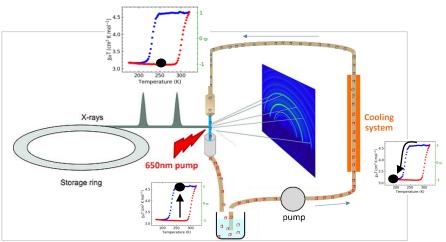


- One pump-probe cycle ↔ diffraction signal usually too low!
- $\rightarrow N$ pump-probe cycles needed for each Δt \rightarrow Study of reversible processes

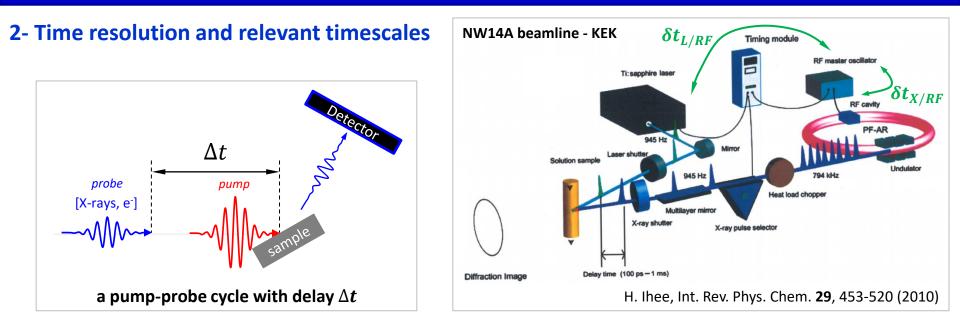
• Irreversible processes : sample must be renewed between each pump-probe sequence ("single-shot" experiments)



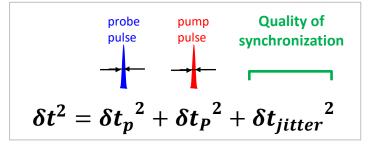
Study of an amorphous-to-liquid-to-liquid photoinduced phase transition. The sample is translated after each pump-probe sequence. [P. Zalden *et al.*, Science **364**, 1062-1067 (2019)]



"Streaming crystallography", Eric Collet M. Hervé et al, Nat Comm **15**, 267 (2024)



Time resolution



Time tools were developed to retrieve δt_{jitter} on a shot-toshot basis and correct the pump-probe delay accordingly. \rightarrow M. Harmand *et al.*, Nature Photonics **7**(3), 215 (2013).

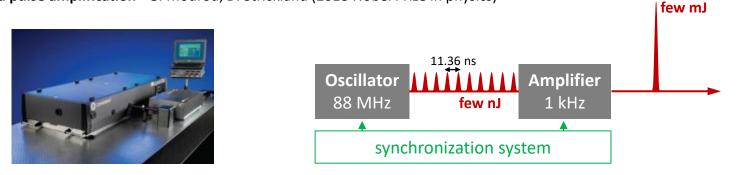
• Timescale of the fastest structural dynamics...

Atomic vibrations: $E_{ph} = \hbar \omega \sim 20 \text{ meV} \Rightarrow T_{osc} \sim 200 \text{ fs}$ Need for < 100 fs time resolution, < 100 fs pump and probe pulses

3a- Sources of ultrashort pump pulses

•Commercially available Ti:Sa lasers (1990 \rightarrow): ~ 40 fs pulses @ 800 nm [1.55 eV], kHz rep rates

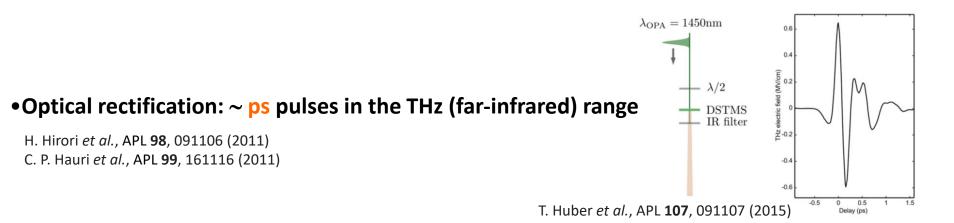
Chirped pulse amplification - G. Mourou, D. Strickland (2018 Nobel Prize in physics)



•Optical parametric amplifiers: \sim 40 fs pulses in the UV to near-IR range

AND LOOK

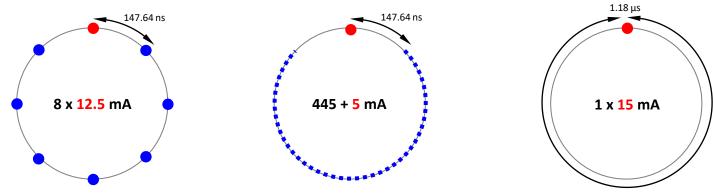
J.-Y. Zhang, "Optical parametric generation and amplification", Routledge; CRC (2019)



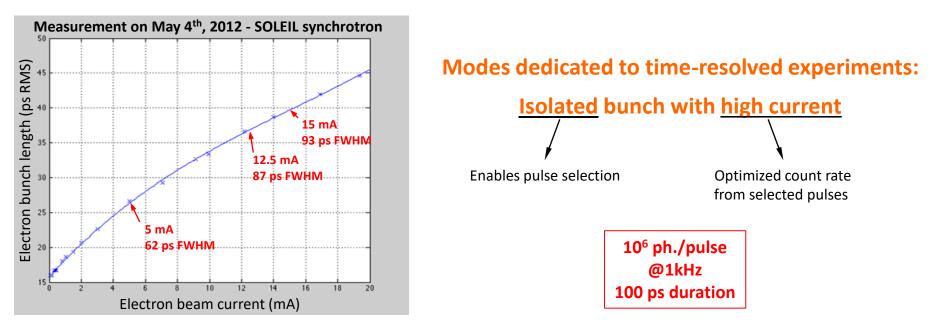
3b- Sources of (ultra)short, hard X-ray pulses [> 5 keV]

<u>1982</u>: first pump-probe diffraction experiment making use of single X-ray pulses (CHESS synchrotron) → B. C. Larson *et al.*, PRL **48**, 337-340 (1982)

100 ps pulses from synchrotrons

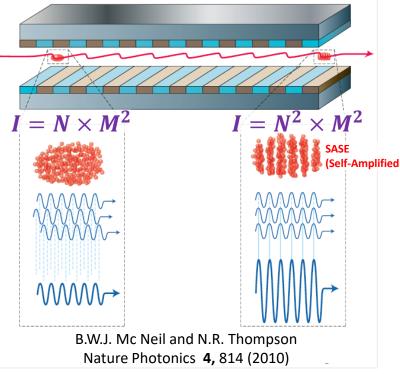


Bunch modes for time-resolved experiments at SOLEIL



3b- Sources of (ultra)short, hard X-ray pulses [> 5 keV]

• few 10 fs X-ray pulses: X-FELs (2009 \rightarrow)





- Short electron pulses produced by a laser-driven electron gun [N electrons]
- Propagation in long undulators (>100 m) [M poles]
- Electron beam bunching

 → Coherent emission of all the electrons

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(Self-Amplified Spontaneous Emission)
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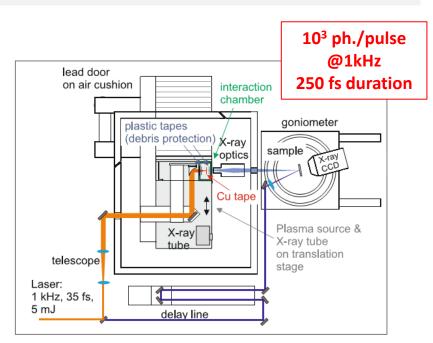
$I \propto N^2 \times M^2$: very high flux 80 fs hard X-ray pulses



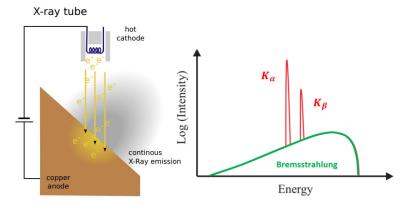
3b- Sources of (ultra)short, Å-wavelength pulses

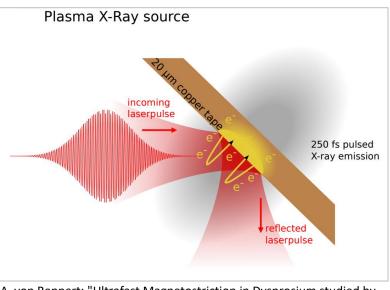
• X-ray plasma sources (1994 \rightarrow)

F. Zamponi, Appl. Phys. A 96, 51-58 (2009)
A. Rousse *et al.*, PRE 50, 2200 (1994)
A. Rousse *et al.*, Nature 410, 65 (2001)



- Laser pulse onto a copper target
- Indirect ionization of Cu atoms
- Emission of X-rays with $K_{\alpha}(Cu)$ wavelength [λ = 1.54 Å]



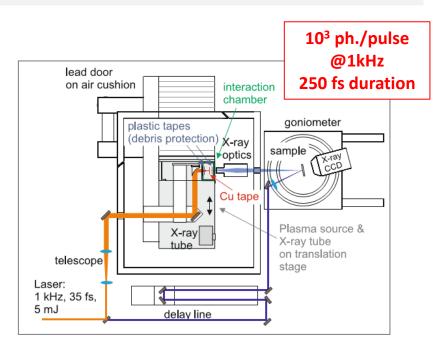


A. von Reppert: "Ultrafast Magnetostriction in Dysprosium studied by Femtosecond X-Ray diffraction" MSc thesis, Universität Potsdam (2015)

3b- Sources of (ultra)short, Å-wavelength pulses

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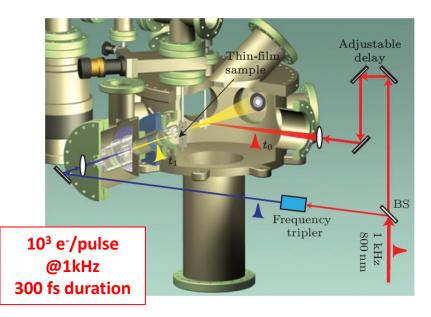
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- Laser pulse onto a copper target
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- Emission of X-rays with $K_{\alpha}(Cu)$ wavelength [λ = 1.54 Å]

• Ultrafast electron diffraction (2003 \rightarrow)

W.-X. Liang *et al.*, Chinese Phys. Lett. **26**, 020701 (2009) R. Srinivasan *et al.*, Helvetica Chimica Acta **86**, 1761-1799 (2003)



- Frequency-tripled Ti:Sa laser pulse ($\lambda = 266$ nm)
- Pulse-driven photocathode \rightarrow photoemission
- Acceleration to \sim 60 keV [$\lambda \sim$ 0.05 Å]



3c- Available pump-probe diffraction setups for users [Europe]

X-ray free electron lasers MID & FXE at Eu-XFEL, Bernina at SwissFEL

• < 100 fs X-ray pulses



~10¹² photons/pulse $\Delta E/E \sim 10^{-3}$ 100 Hz [SwissFEL] 10 Hz/4.5 MHz [Eu-XFEL]

Short pulse facility, Lund (Sweden) femtoMAX beamline

• < 160 fs X-ray pulses



10⁶ photons/pulse $\Delta E/E \sim 10^{-2}$

10 Hz

Synchrotrons ID09 at ESRF, KMC3-XPP at BESSYII

• 100 ps X-ray pulses

• 10 ps X-ray pulses (low-α, BESSY)



10⁶ photons/pulse $\Delta E/E \sim 10^{-4}$ 1 Hz to 3 kHz [ID09] 1 to 625 kHz [XPP]

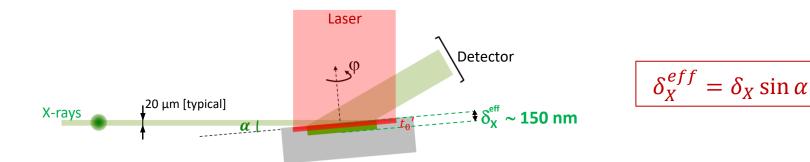
~10³ photons/pulse $\Delta E/E \sim 10^{-4}$

1 to 625 kHz [XPP]

4a- Methods specific to pump-probe X-ray diffraction : grazing incidence geometry



- Nanocrystals (dilute) and thin-films can be studied using unconstrained diffraction geometries...
- Bulk crystals: a grazing incidence of the X-ray beam is required



Beware - The grazing angle α should be small, but larger that the critical angle for total external reflection α_c :

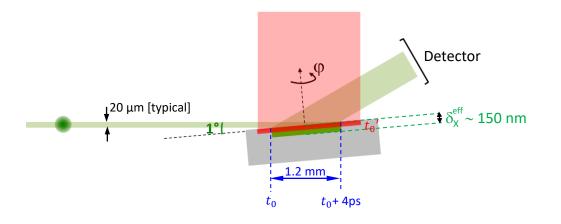
$$\alpha_c \approx \lambda \sqrt{rac{
ho_{at} r_e}{\pi}}$$

 r_e : classical radius of the electron (2.818 10^{-15} m)

Orders of magnitude at 7 keV: $\alpha_c \sim 0.3 - 0.4^\circ \implies \alpha \sim 0.5 - 1^\circ$

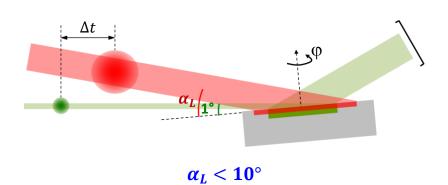
4b- Methods specific to pump-probe X-ray diffraction : collinear X-ray and laser beams

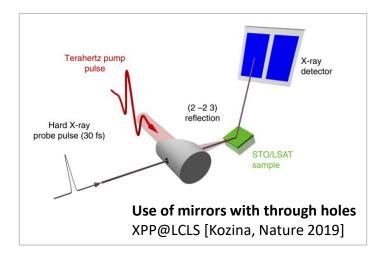
• Grazing incidence geometry



!! Loss of effective time resolution from the difference in relative arrival times between the pump and X-ray beams !!

• Collinear pump & probe beams





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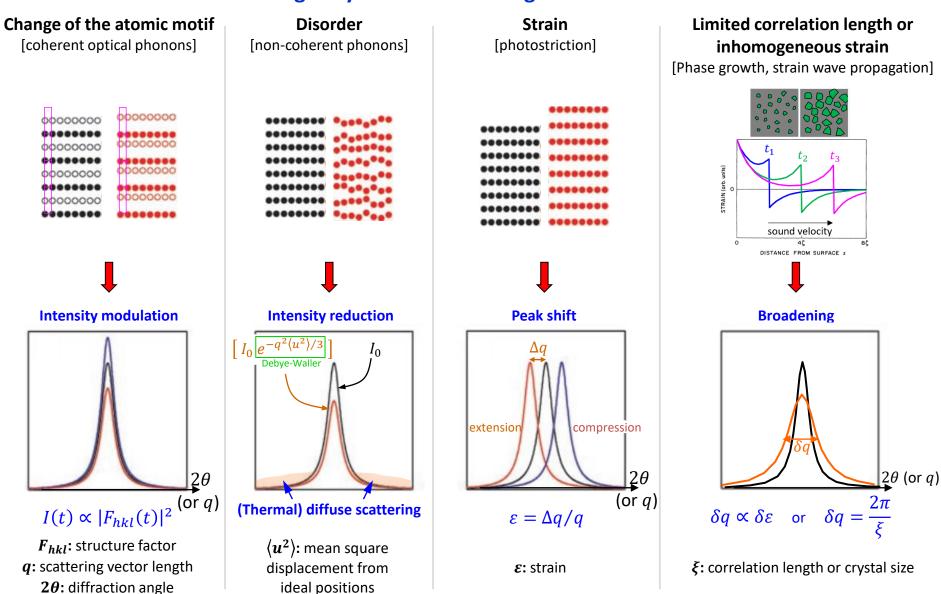
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- Coherent phonons
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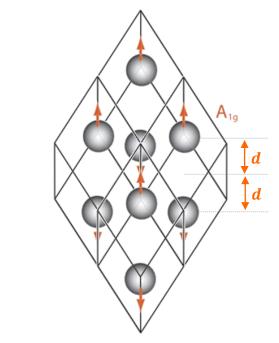
1- Structural information brought by the diffraction signals

M. Bargheer et al., ChemPhysChem 7, 783 – 792 (2006)



2- Coherent optical phonons

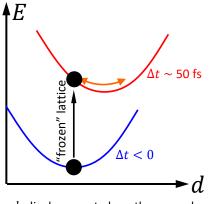
<u>Coherent phonons</u> give rise to atomic motions with a fixed phase relation between them, homogeneously in the sample



Example: coherent A1g phonon in bismuth Atoms are moving in-phase from one unit cell to another

• Displacive excitation of coherent phonons (DECP)

Strong coupling between the initial/final electronic state and one particular vibrational mode \rightarrow Shift of the local minimum of the potential energy along the vibrational mode coordinate

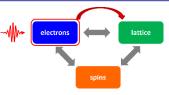


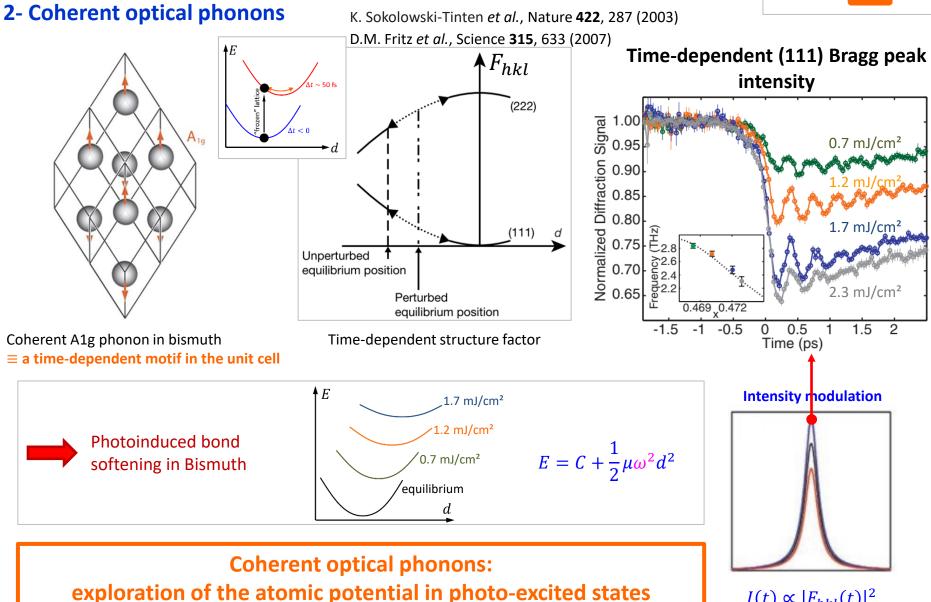
d: displacement along the normal coordinates of the A_{1g} mode

<u>All</u> the atoms of the excited volume are simultaneously significantly displaced with respect to the new quasi-equilibrium position: they start to oscillate in phase along the normal coordinate of the mode.

Further reading on coherent optical phonons:

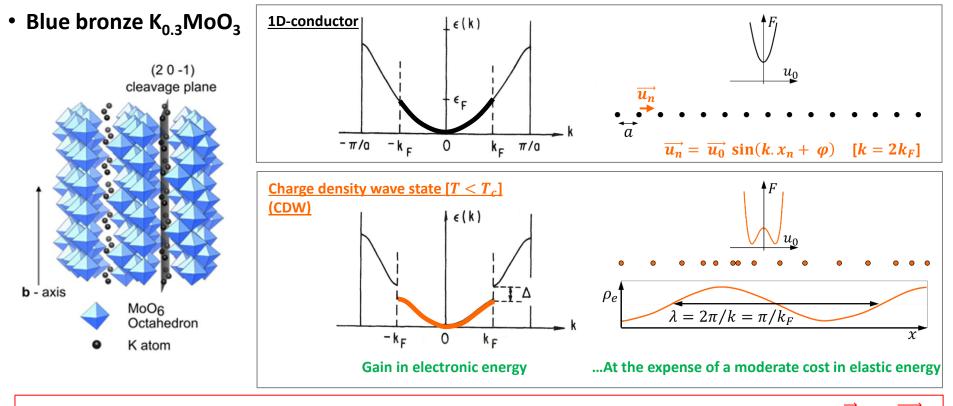
R. Merlin, Solid State Commun. **102**, 207-220 (1997) H.J. Zeiger *et al.*, Phys. Rev. B **45**, 768-778 (1992) Y.-X. Yan *et al.*, J. Chem. Phys. **83**, 5391 (1985) T. Stevens *et al.*, Phys. Rev. B **65**, 144304 (2002)





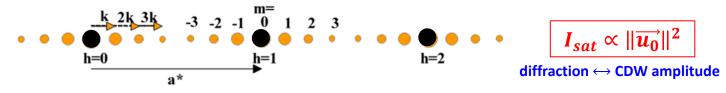
 $I(t) \propto |F_{hkl}(t)|^2$

3a- Photoinduced phase transitions (PIPTs) achieved through coherent atomic motions



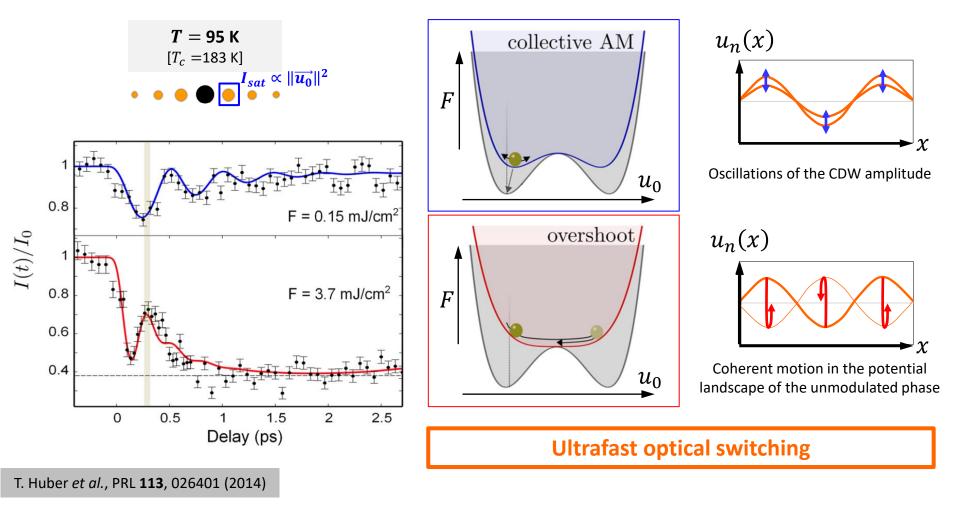
Strong e-ph coupling, CDW transition occurs along the coordinate u_0 of the vib. mode with $ec{k}=2ec{k_F}$

• Diffraction pattern: satellite peaks



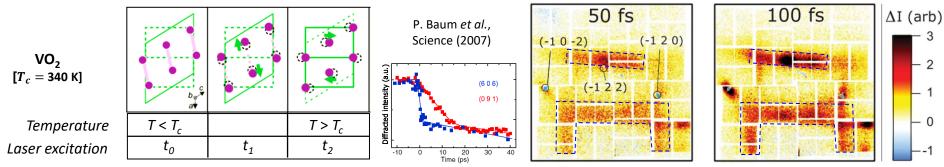
3a- Photoinduced phase transitions (PIPTs) achieved through coherent atomic motions

Ultrafast PIPT in Blue bronze K_{0.3}MoO₃



See also [manganite Pr_{0.5}Ca_{0.5}MnO₃]: P. Beaud *et al.*, Nature Materials **13**, 923 (2014)

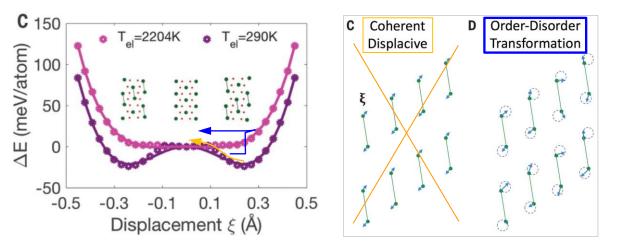
3b- Photoinduced phase transitions (PIPTs): non-fully coherent atomic motions !

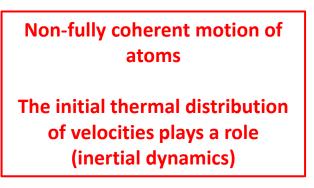


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Detector images: I(t) - I_0
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Insulator to metal transition revisited using an X-FEL (LCLS, Stanford):

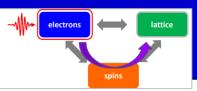
- \rightarrow Drop of diffracted intensity of the $\overline{1}0\overline{2}$, $\overline{1}20$ and $\overline{1}22$ reflections on the 100 fs timescale
- \rightarrow Developement of a diffuse scattering signal on the 100 fs timescale





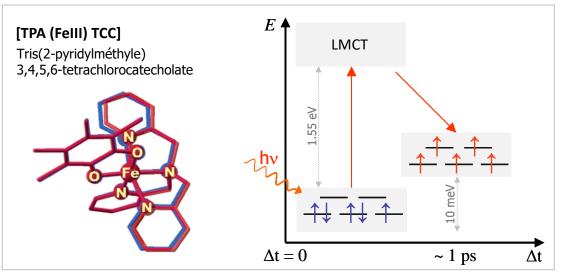
spins

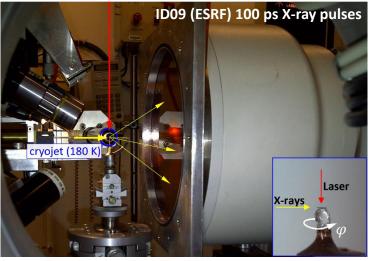
S. Wall et al., Science 362, 572-576 (2018)

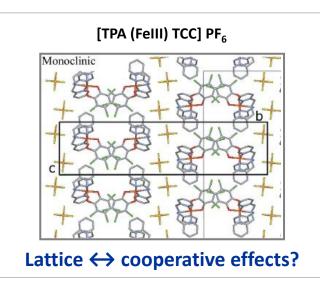


4- Photoinduced strain and temperature elevation (picosecond to microsecond timescales)

Spin state switching in a molecular crystal

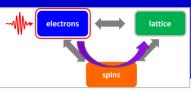






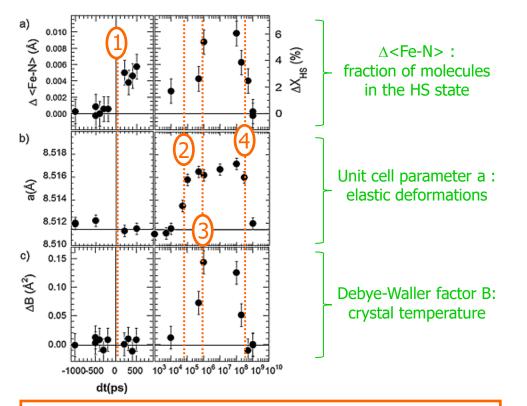
- 360° rotation in $oldsymbol{arphi}$: ~ 3300 Bragg peaks
- → **Structural refinement**, starting from the structure observed at equilibrium
- \bullet Crystal thickness (15 μm) is chosen roughly equal to the penetration depth of 800 nm laser pulses
- \rightarrow Near-homogeneous excitation
- Circularly polarized laser beam
- \rightarrow Constant absorption probability of the 800 nm photons during φ rotations

M. Lorenc et al., PRL 103, 028301 (2009)



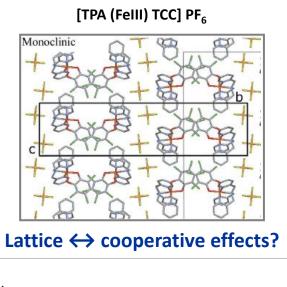
4- Photoinduced strain and temperature elevation (picosecond to microsecond timescales)

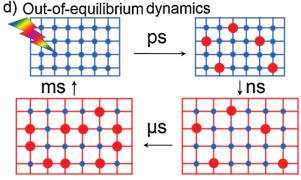
Spin state switching in a molecular crystal



Photoinduced spin state switching in [TPA Fe(III) TPP] PF₆:

- 1) Laser-induced switching of molecules
- 2) Strain wave propagation
- 3) Heat diffusion + additional spin state switching
- 4) Relaxation





M. Lorenc et al., PRL 103, 028301 (2009)

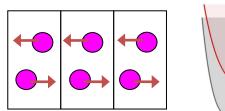
Switching due to strain waves (nanocrystals): R. Bertoni *et al.*, Nat. Mater. **15**, 606 (2016)

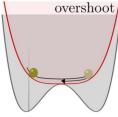
Photoinduced structural dynamics

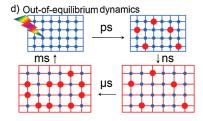
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Out-of-equilibrium structural dynamics:

- Coherent phonons
- Photoinduced phase transitions
- Strain wave propagation
- Heat diffusion





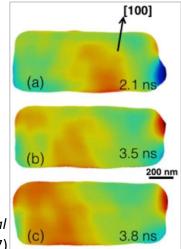


Pump-probe diffraction

- Bragg peak intensities: time-dependent motif, volume fraction, average temperature (Debye-Waller factor)
- Bragg peak positions: transient strain
- Bragg peak widths: correlation length (up to 100 nm)
- Diffuse scattering: short range correlated atomic displacements

Pump-probe diffraction at X-FELs:

 $\bullet\,\mu\text{-diffraction,}$ coherent diffractive imaging, resonant scattering



Strain pulse in a ZnO nanocrystal M. J. Cherukara, Nano Lett. **17**, 1102-1108 (2017)