# a view from time-resolved X-ray diffraction





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#### I. Scientific motivations

#### II. Pump-probe diffraction

- Principle
- Time resolution & synchronization
- Short X-ray pulse sources
- Specific geometrical constraints

#### III. Example

- Photo-induced phase transition in K<sub>0.3</sub>MoO<sub>3</sub>



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## **Structural dynamics in physics**

#### • Crystals at thermodynamic equilibrium

Atomic displacements : sum of normal modes  $\overrightarrow{u_n}(\vec{r},t) = \sum_{\lambda, \|\vec{k}\|} u_n(\lambda, \vec{k}) \vec{e}_{\lambda, \vec{k}} e^{i[\omega(\lambda, \vec{k})t - \vec{k}.\vec{r}]}$ 



Transverse mode

→ Experiments in the <u>frequency domain</u>: inelastic neutron scattering, Raman scattering...

#### Photo-induced structural dynamics



→ Experiments in the time domain: time-resolved pump-probe diffraction



## Structural dynamics in physics (1/2)





Exploration of the potential in photo-excited states

#### Novel states of matter Ultrafast control of the physical properties



## Structural dynamics in physics (2/2)





**Cooperative effects in molecular crystals** Light control of ferroelectric and magnetic materials

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## Time-resolved pump-probe diffraction

#### Following photo-induced structural changes as a function of time



- One pump-probe cycle  $\leftrightarrow$  diffraction signal too low !
- $\rightarrow N$  pump-probe cycles needed for each  $\Delta t$ : study of reversible processes
- ightarrow Irreversible processes : liquid jets / serial crystallography at X-FELs





I. Schlichting, IUCrJ **2**, 246–255 (2015) V. Panneels *et al.*, Structural Dynamics **2**, 041718 (2015) T. R. M. Barends, Science **350**, 445 (2015)

## Typical layout for pump-probe diffraction



•Commercially available Ti:Sa lasers (1990  $\rightarrow$ ): ~ 40 fs pulses @ 800 nm [1.55 eV]











## X-ray pulse sources

#### • Few ps pulses from synchrotrons: low- $\alpha$ mode

#### Normal operation:

- Optics optimized for a low-emittance electron beam
- Dispersion of  $E_{e^-} \Longrightarrow$  dispersion of  $e^-$  revolution period
- Elongated e<sup>-</sup> bunches, longer X-ray pulses



 $\rightarrow$  User operation at BESSY, SOLEIL, DIAMOND



## X-ray pulse sources

• 100 fs X-ray pulses from synchrotrons: femto-slicing sources (1995  $\rightarrow$ )



• 
$$\vec{E} \cdot \vec{v} \neq 0$$

# $\Lambda_w = \Lambda_L$

#### **Co-propagation of an electron bunch** and a laser pulse in a wiggler

 $\rightarrow$  Modulation of electron energies in the overlap zone



A - 1 L.

#### → BESSY (soft X-rays), ALS, SLS & SOLEIL (hard X-rays)

## X-ray pulse sources

• 100 fs X-ray pulses: X-FELs (2009  $\rightarrow$ )



- Short electron pulses produced by a laser-driven electron gun [N electrons]
- Propagation in long undulators (100 m) [M poles]
- Electron beam bunching
   → Coherent emission of all the electrons

### $I \propto N^2 \times M^2$ : very high flux 80 fs hard X-ray pulses





## Time-resolved pump-probe diffraction: laser-based sources

• X-ray plasma sources (1994  $\rightarrow$ )



- Laser pulse onto a copper target
- Indirect ionization of Cu atoms
- Emission of X-rays with  $K_{\alpha}(Cu)$  wavelength [ $\lambda$ = 1.54 Å]

10<sup>3</sup> ph./pulse @1kHz 100 fs duration

F. Zamponi, Appl. Phys. A 96, 51-58 (2009)
A. Rousse *et al.*, PRE 50, 2200 (1994)
A. Rousse *et al.*, Nature 410, 65 (2001)

• Ultrafast electron diffraction (2003  $\rightarrow$ )



- Frequency-tripled Ti:Sa laser pulse ( $\lambda = 266 \text{ nm}$ )
- Pulse-driven photocathode  $\rightarrow$  photoemission
- Acceleration to  $\sim$  60 keV [ $\lambda \sim$  0.05 Å]

#### 10<sup>3</sup> e<sup>-</sup>/pulse @1kHz 300 fs duration

W.-X. Liang *et al.*, Chinese Phys. Lett. **26**, 020701 (2009) R. Srinivasan *et al.*, Helvetica Chimica Acta **86**, 1761-1799 (2003)

## Pump-probe diffraction : typical photon or electron fluxes

Synchrotrons [repetition rate 1 kHz]			
• 80 ps X-ray pulses	$\bigcirc\bigcirc\bigcirc\bigcirc$	<b>10<sup>6</sup> photons/pulse</b> $\Delta E/E \sim 10^{-4}$	10 <sup>9</sup> ph/s
• Few ps X-ray pulses (low- $\alpha$ )		<b>10<sup>4</sup> photons/pulse</b> $\Delta E/E \sim 10^{-3}$	10 <sup>7</sup> ph/s
• 100 fs X-ray pulses (femto-slicing)	No. of the second secon	<b>10<sup>3</sup> photons/pulse</b> $\Delta E/E \sim 8.10^{-3}$	10 <sup>6</sup> ph/s
Laser-based sources [repetition rate 1 kHz]			
• 100 fs X-ray pulses (plasma source)		10 <sup>3</sup> photons/pulse $\Delta E/E \sim 10^{-4}$	10 <sup>6</sup> ph/s
• 300 fs electron pulses		10 <sup>3</sup> electrons/pulse	10 <sup>6</sup> e <sup>-</sup> /s
X-ray free ele	ctron lasers [repe	tition rate 100 Hz]	
• 80 fs X-ray pulses		<b>10<sup>11</sup> photons/pulse</b> $\Delta E/E \sim 10^{-3}$	10 <sup>13</sup> ph/s
		ATA A	



## **Time resolved X-ray diffraction: experimental facts**

• X-rays and IR photons: differing penetration depths !



[Typical values for hard condensed matter, 7 keV X-ray photons ]



• Grazing incidence geometry, pump & probe beams collinear





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## Atomic structure of blue bronze $(K_{0.3}MoO_3)$

• A quasi-one-dimensional conductor...



 $\sigma_b
 = 3 \ 10^2 \ [\Omega.cm]^{-1}$ 

...which undergoes a transition to a charge density wave phase at 183 K



J. Graham and A.D. Wadsley, Acta Cryst. **20**, 93 (1966) G. Grüner, "Density waves in solids"

## Formation of a charge density wave - Peierls model

• A metal-insulator transition driven by a periodic lattice distortion





Appearance of a charge density wave in blue bronze  $(K_{0.3}MoO_3)$ 

• K<sub>0.3</sub>MoO<sub>3</sub>: satellite peaks @  $(h k l) + (1 q_b \frac{1}{2})$ 





## Ultrafast light control of the physical properties of CDW compounds ?



• Femto-slicing source @ SLS:



• Time-dependence of the satellite  $\left(1 \left[4-q_b\right] \frac{1}{2}\right)$  - Low fluence



• Time-dependence of the satellite  $\left(1 \left[4-q_b\right] \frac{\overline{1}}{2}\right)$  - Higher fluences



•  $F = 1 \text{mJ/cm}^2$ 

- $\rightarrow$  The recovery time of satellite peak intensity increases
- $\rightarrow$  Coherent oscillations: hardly observable

#### • $F > 1 \text{mJ/cm}^2$

- $\rightarrow$  No recovery of satellite peak intensity within 10 ps
- $\rightarrow$  Oscillation frequency doubled w/r to the low fluence case

#### Significant changes of the atomic potential surface

T. Huber et al., PRL 113, 026401 (2014)

• Time-dependence of the satellite  $\left(1 \left[4-q_b\right] \frac{\overline{1}}{2}\right)$  - Higher fluences



• Free energy vs laser excitation  $[\eta \propto \text{laser fluence}]$ 

$$F = F_0 + \frac{1}{2} \left[ \eta e^{-\frac{t}{\tau}} - 1 \right] \, u_0^2 + \frac{1}{4} {u_0}^4$$

#### Equation of motion to be solved:

$$C_1 \frac{\partial^2 u_0(t)}{\partial t^2} = -\frac{dF}{du_0} - C_2 \gamma(t) \frac{\partial u_0(t)}{\partial t}$$

Non-harmonic motions of atoms

T. Huber et al., PRL 113, 026401 (2014)

• Time-dependence of the satellite  $\left(1 \left[4-q_b\right] \frac{\overline{1}}{2}\right)$  - Higher fluences



Ultrafast change of atomic potential symmetry

# Thank you !

